



POLITECNICO DI MILANO

Piazza Leonardo da Vinci, 32 - 20133 Milano
Tel. +39.02.2399.1 - <http://www.polimi.it>



Advanced course on

**HIGH RESOLUTION ELECTRONIC MEASUREMENTS
IN NANO-BIO SCIENCE**

INSTRUMENTATION FOR NOISE MEASUREMENTS

Noise as Signal

Marco Sampietro

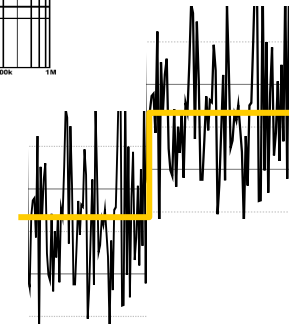
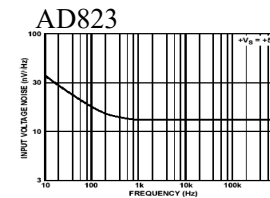
Outline of the lesson

We have been FIGHTING all along the course AGAINST NOISE !!

Noise is universal and unavoidable (thermal noise, shot noise, $1/f$ noise, ...)

In electronic devices & circuits it is important to measure the noise ...

... as it sets the minimum detectable signal.



Can NOISE be OUR FRIEND ?

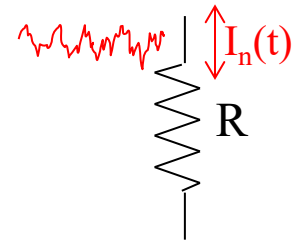
Use it as a macroscopic «echo» of the micro-physics ruling the device properties

How to MEASURE very low NOISE ?

What is the noise of a resistor?

For **any resistor** (dissipative system)
in thermal equilibrium :

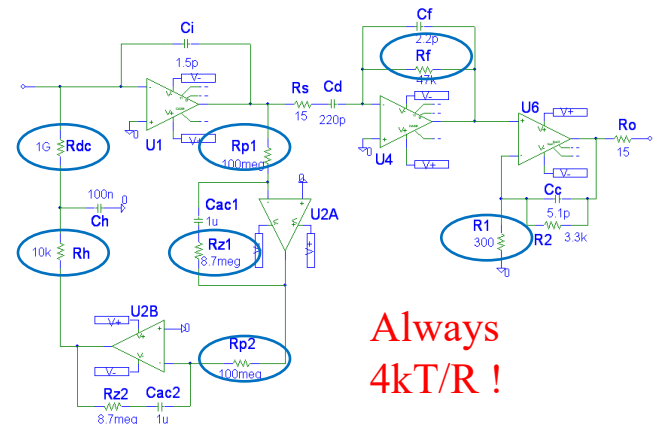
$$S_i = \frac{4kT}{R}$$



The thermal noise is commonly used
also in non-equilibrium conditions

Why ?

Inelastic scattering, by dissipating the energy that electrons gain from the field and randomizing the momentum, reduces the electron average energy to that of the lattice, and therefore the noise corresponds to thermal-Nyquist noise at any bias.



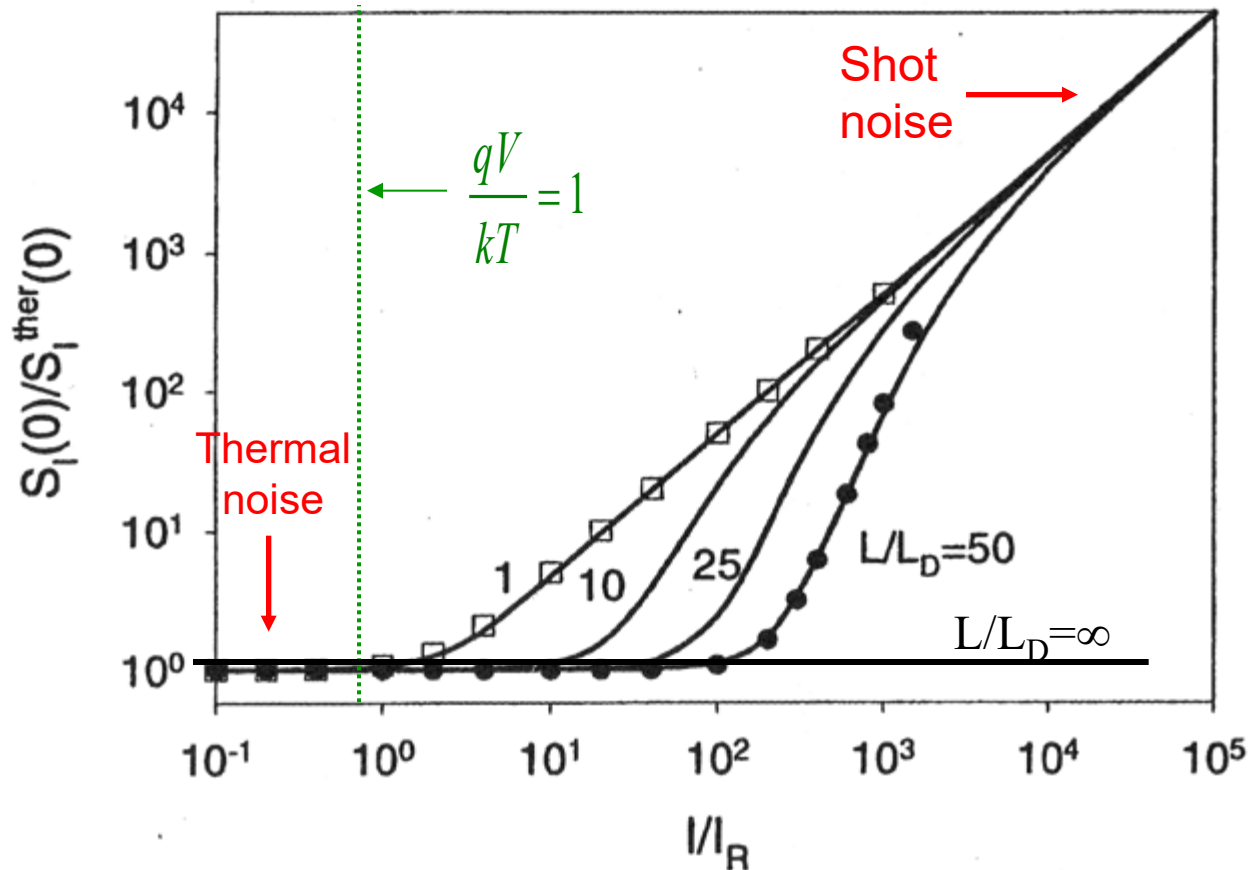
Always
 $4kT/R$!

Why we do not see any granularity of the charge (shot noise) ?

$$S_i = 2qI \quad [A^2/Hz]$$

The prediction of the G-R theory

G.Gomila, L.Reggiani, *Phys. Rev. B* 62 (2000), p.8068



Departure from the thermal noise:

$$I_1 = \left(\frac{L}{L_D} \right)^{4/3} I_R \quad I_R = \frac{V_{th}}{R}$$

($L > L_D$; per $L < L_D \rightarrow I_1 = I_R$)

Shot noise:

$$I_2 = \left(\frac{L}{L_D} \right)^2 I_R \quad (\tau_T = \tau_d)$$

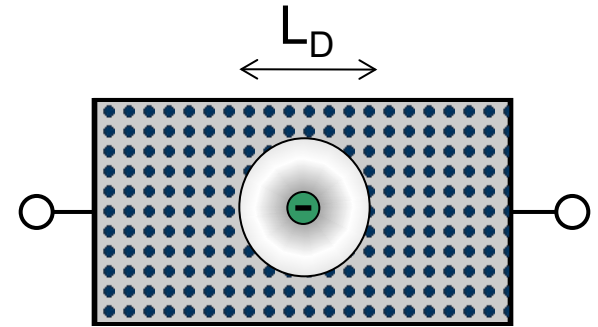
($L > L_D$; per $L < L_D \rightarrow I_2 = I_R$)

A macroscopic resistor might show shot noise !

Key parameters

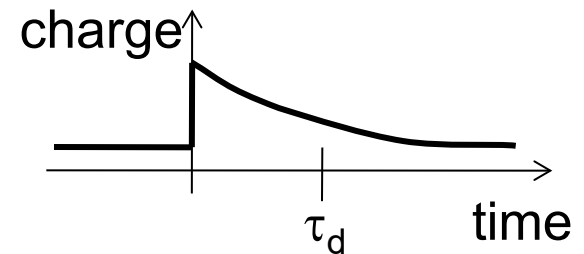
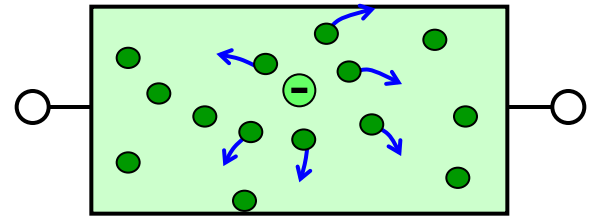
$$L_D = \sqrt{\frac{\varepsilon V_{th}}{q n}}$$

Debye length



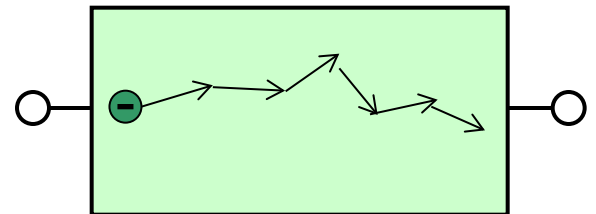
$$\tau_d = \frac{\varepsilon}{q n \mu}$$

Dielectric relaxation time



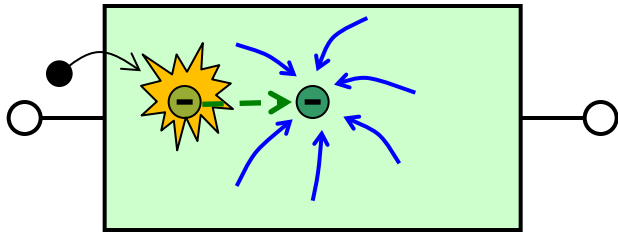
$$\tau_T = \frac{L^2}{\mu V}$$

Transit time



Thermal noise in resistors (“standard” case)

transit time $\gg \tau_d$ (dielectric relaxation time)

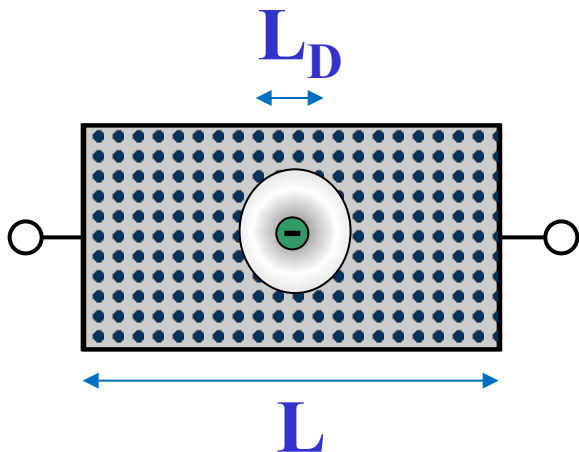


Time long enough to shield the carrier
→ electrodes “do not see” the single carrier
but a “collective” effect



Thermal noise

Alternatively :

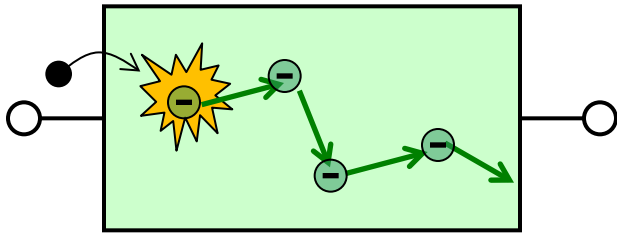


$$L \gg L_D$$

The carrier is shielded
→ electrodes “**do not see**” the single carrier

Shot noise in resistors

transit time $\ll \tau_d$ (dielectric relaxation time)



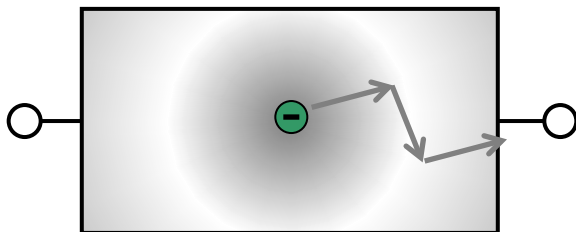
The material has no time to shield the carrier \rightarrow electrodes “see” the single charge

(mean time btwn collisions) $\tau_m \ll$ transit time



Shot noise

$L_D \gg L$



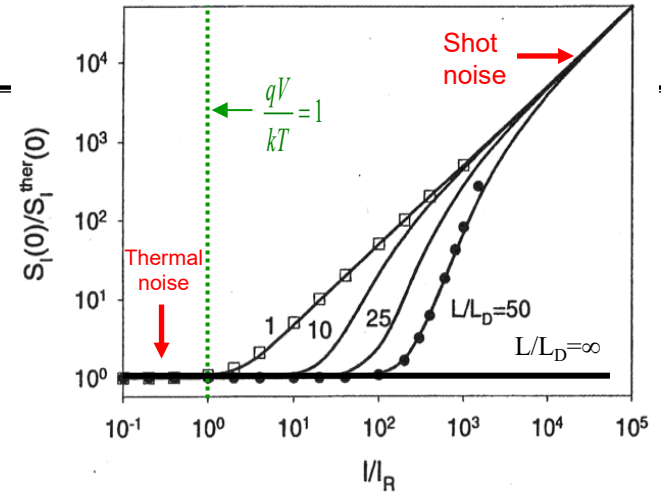
Charge carriers are independent.
 \rightarrow electrodes “**see**” the single carrier.

Fluctuation of the number of carrier due to random scattering

Independent of the injection of carriers: it is the random motion in the material that gives the shot noise

(carriers thermalize but do not correlate)

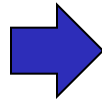
Examples



- “Standard” resistor:

$R=100\text{k}\Omega$

$L=1\text{mm}$



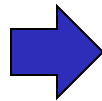
$L_D=0.5\text{nm}$

Thermal-shot transition at $E=10\text{ GV/mm} !!!$

- Heavily doped silicon resistor:

$n=10^{17}\text{ cm}^{-3}$

$L=1\text{mm}$



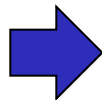
$L_D=12.5\text{nm}$

Thermal-shot transition at $E=9\text{ kV/mm}$ ($8\text{V}/\mu\text{m}$)

- Lightly doped silicon resistor:

$n=10^{14}\text{ cm}^{-3}$

$L=1\mu\text{m}$



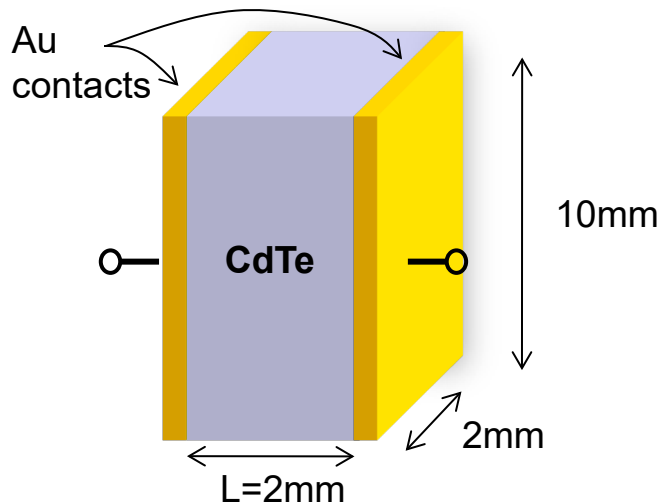
$L_D=400\text{nm}$

Shot noise at $E=1.6\text{ kV/cm}$ ($160\text{mV}/\mu\text{m}$)
(transition at 850 V/cm)

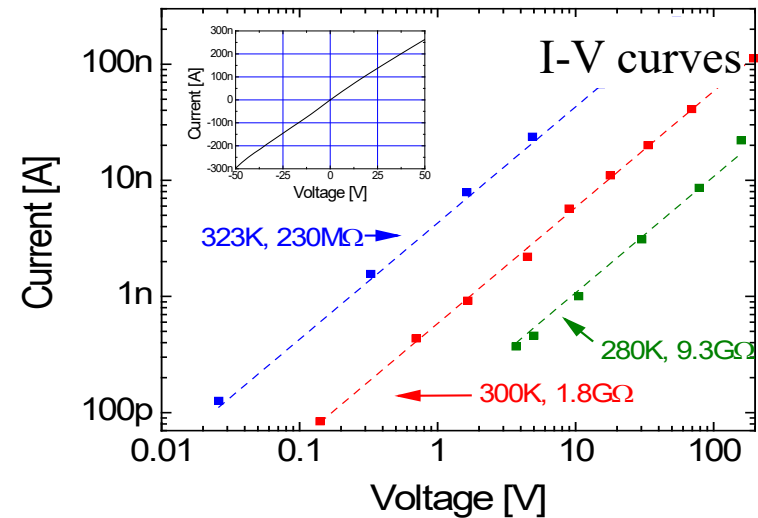
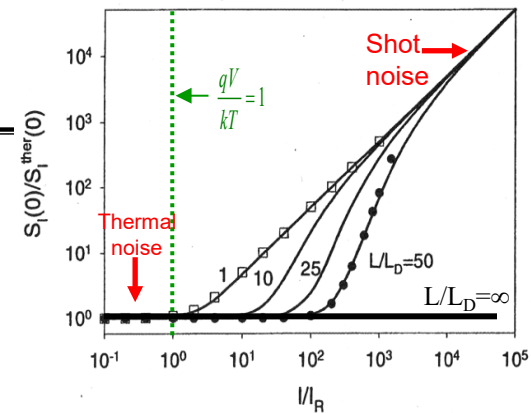
Experimental validation on CdTe

CdTe crystal (ohmic contact in gold):

- Wide band gap (1.47eV)
- Lightly doped ($p \cong 9 \cdot 10^7 \text{ cm}^{-3} |_{T=300\text{K}}$)
- $\rho = 1.8 \text{ G}\Omega \cdot \text{cm}$ ($L/L_D = 4.4$)
- Mobility is electric field independent up to tens of kV/cm

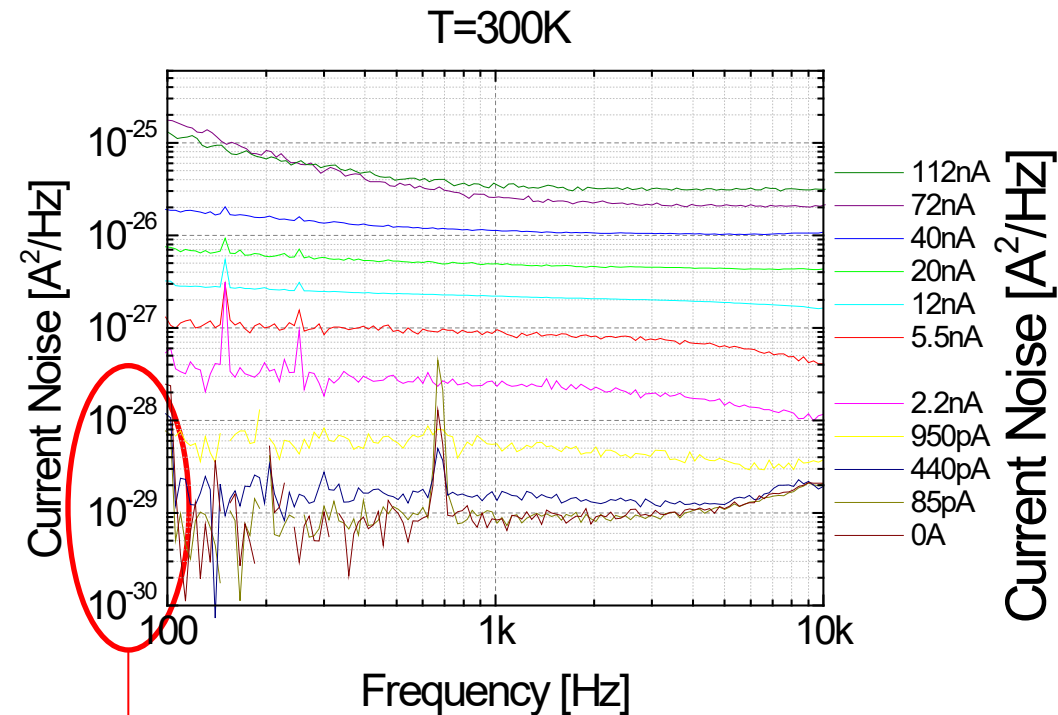


Macroscopic!

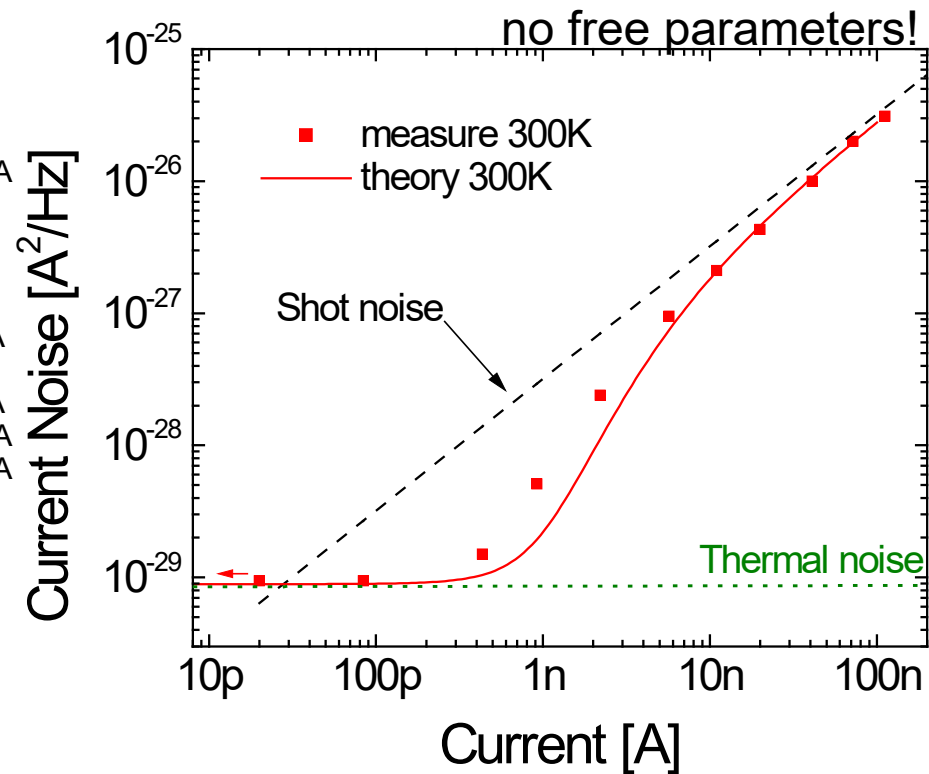


Good linearity (and symmetry) at different temperatures
Exponential dependence on the temperature ($n = n_0 \exp(-E/kT)$)

Shot noise of a resistor (at room temperature)



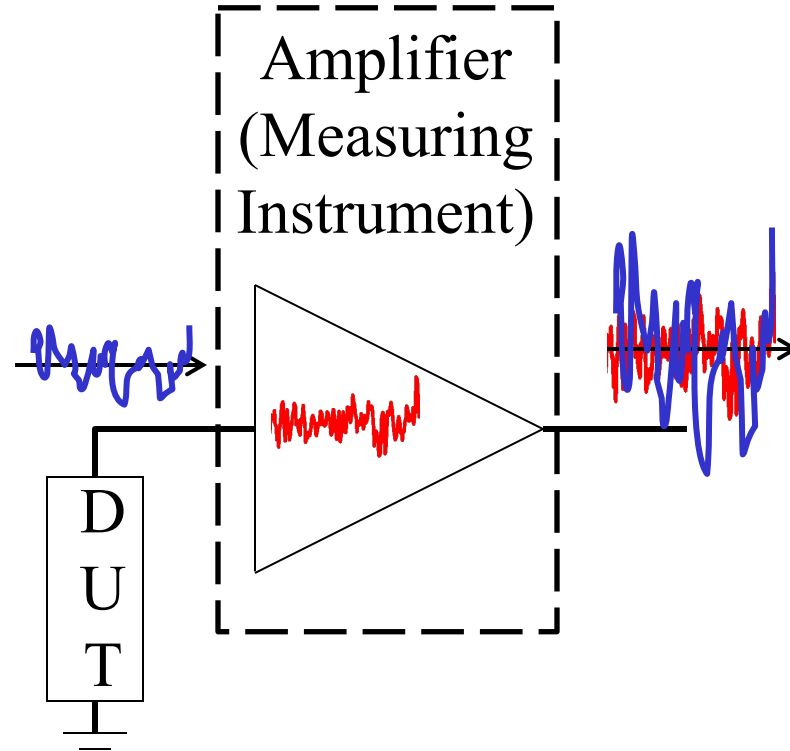
Very small values of noise !



G. Ferrari et al., APL, 83, 2450 (2003)
G. Gomila et al., PRL, 92, 226601 (2004)

T [K]	p [cm^{-3}]	L/L _D
280.5	2×10^7	2
300	9×10^7	4.4
323	7×10^8	11.8

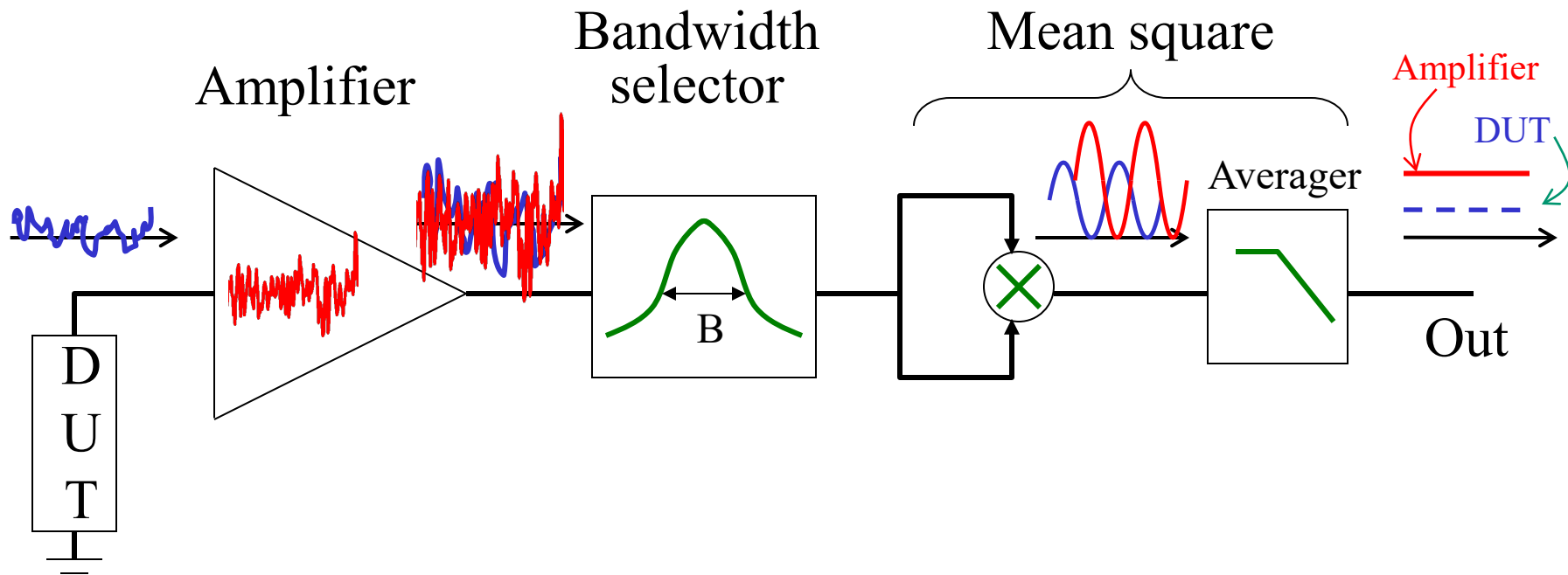
How to measure very low noise ?



Noise power of input amplifier **adds** to the DUT signal power and therefore sets the minimum detectable “DUT NOISE” signal.

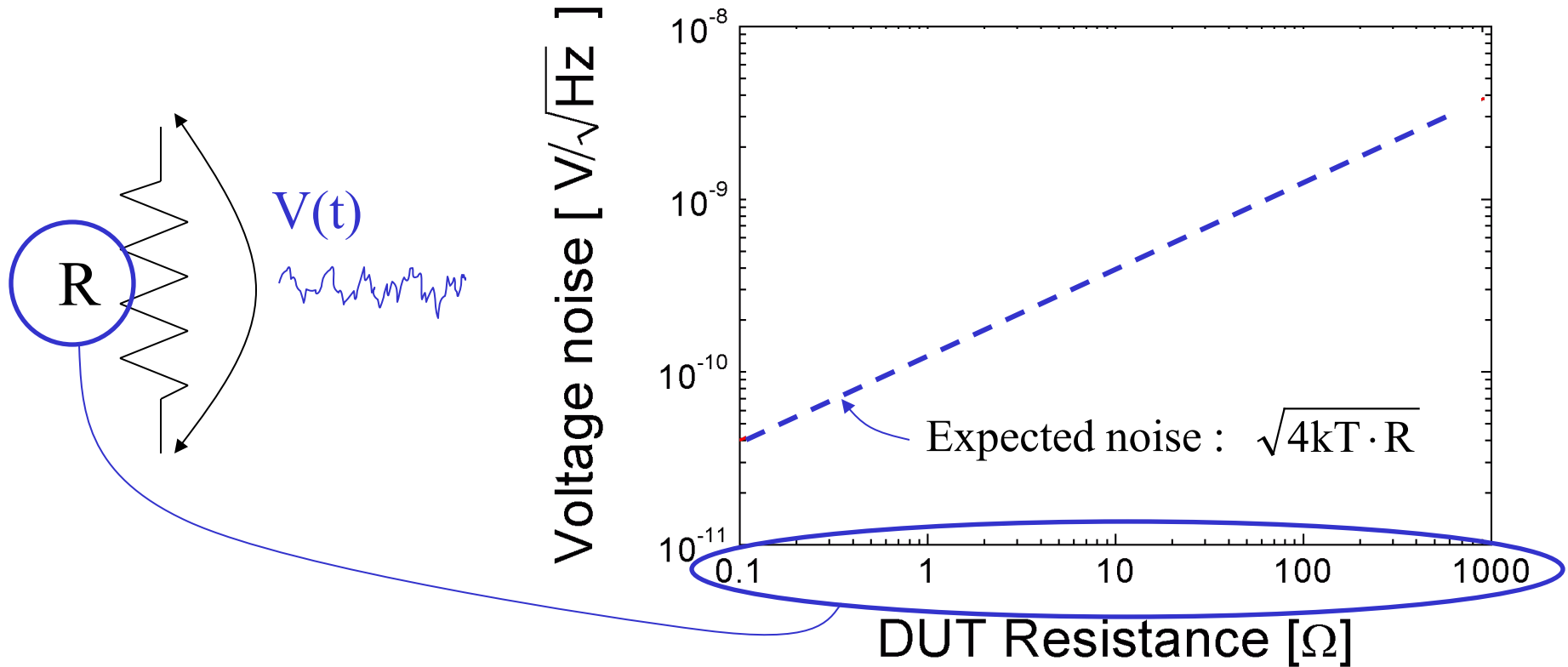
This apply frequency by frequency

A standard spectrum analyzer

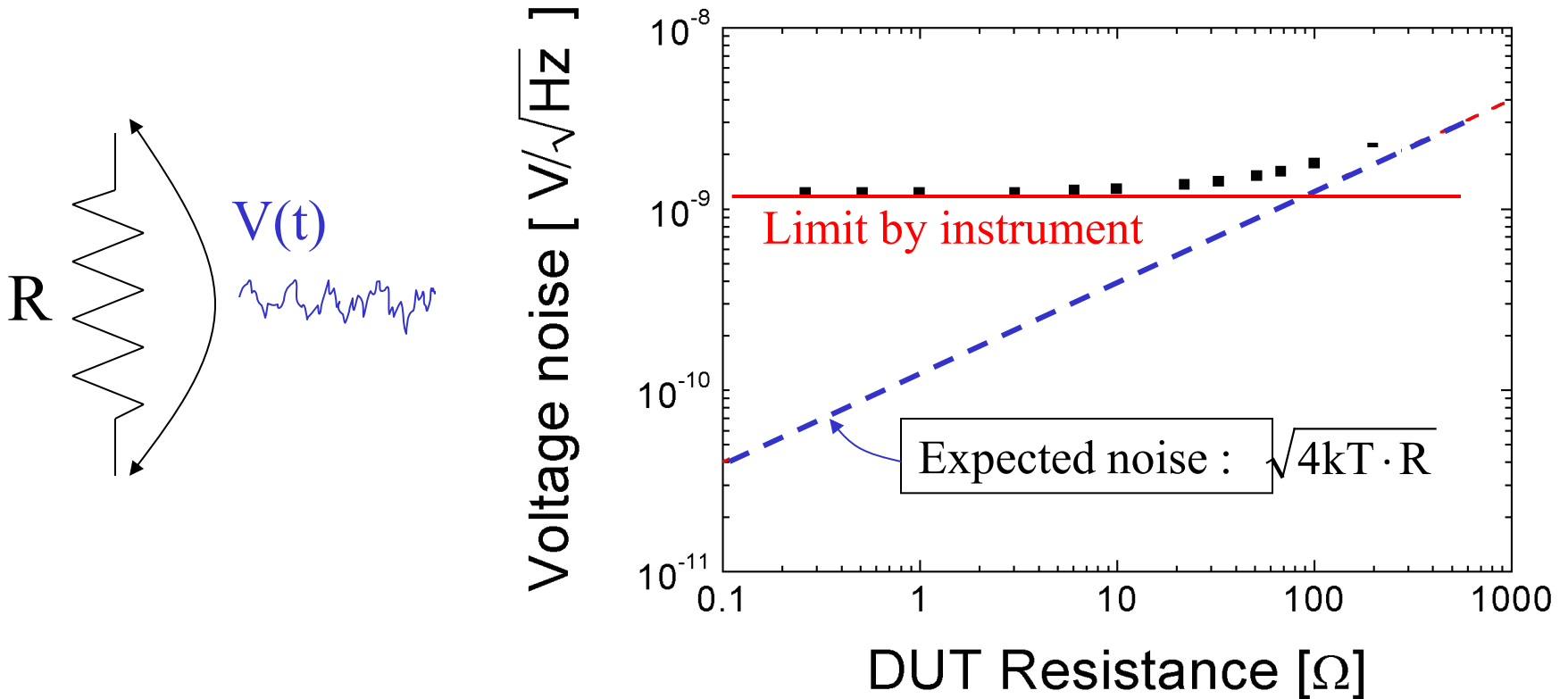


Typical white noise power of best commercial analyzers is $1\text{nV}/\sqrt{\text{Hz}}$.
($10^{-18} \text{ V}^2/\text{Hz}$)

An example

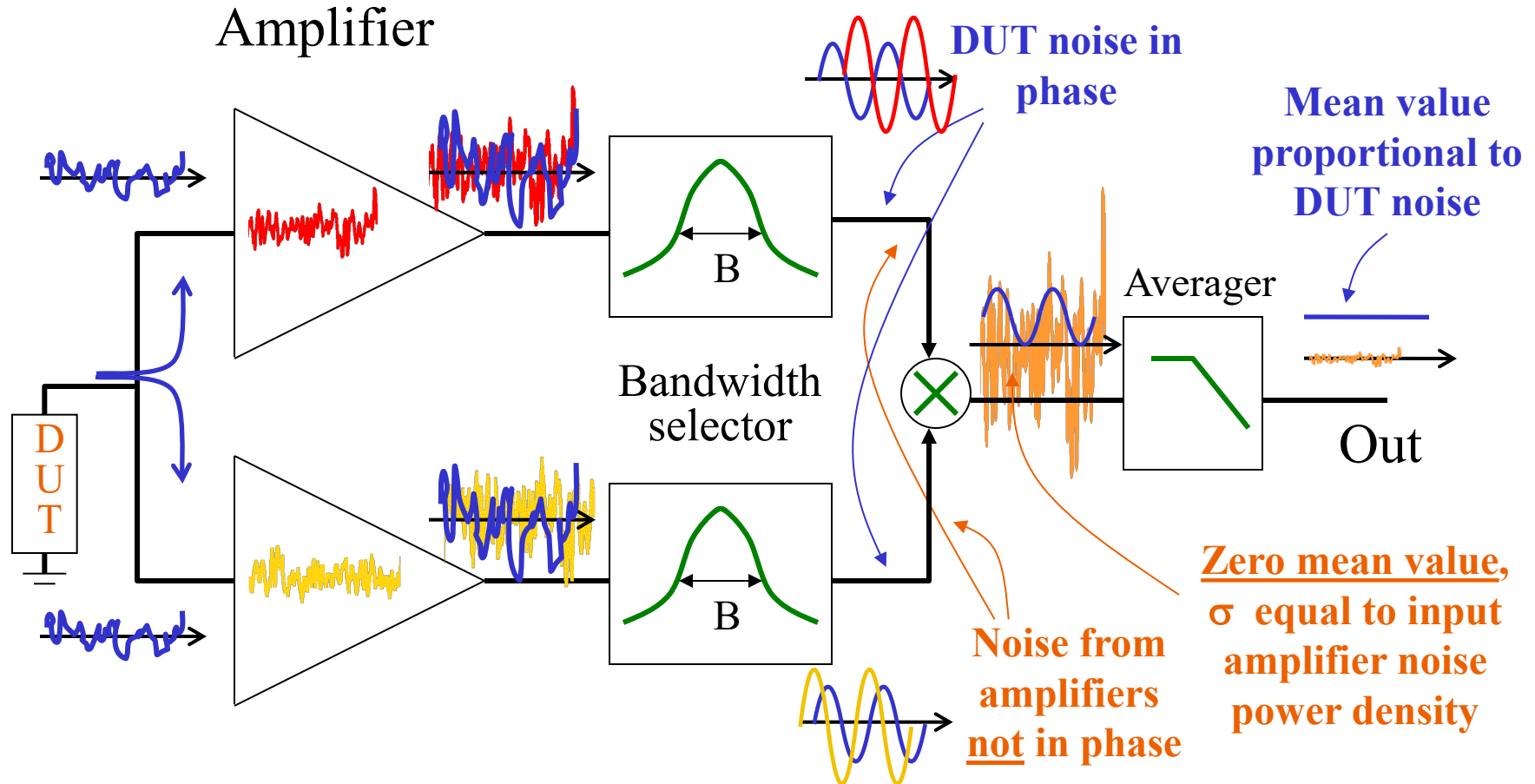


An example



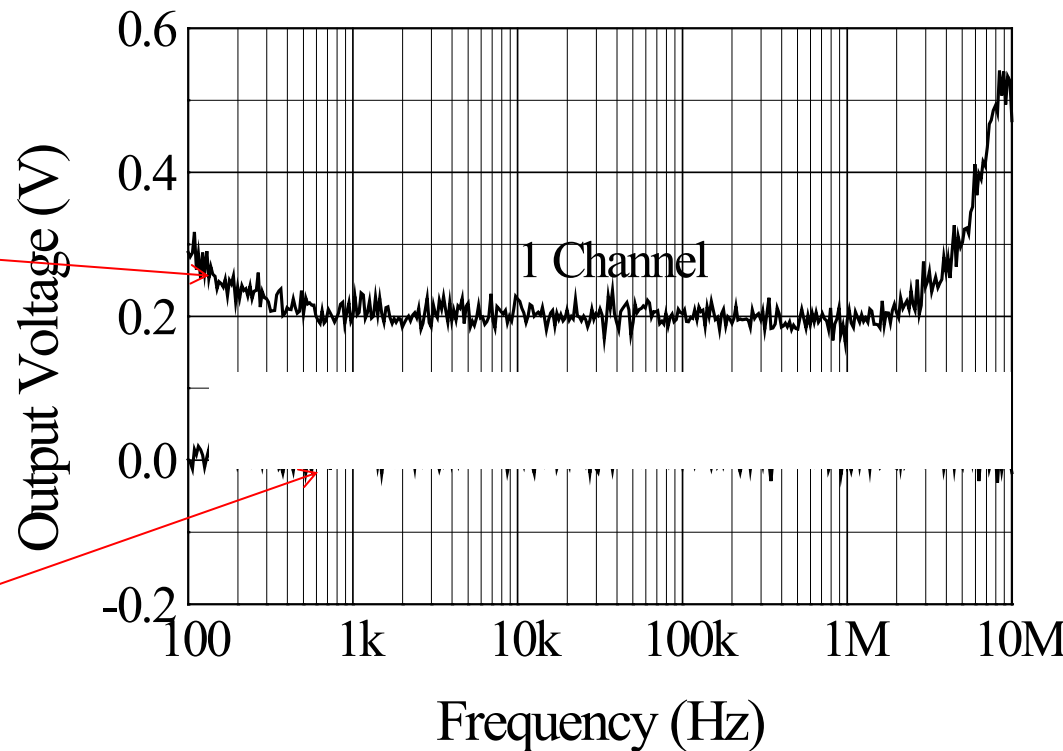
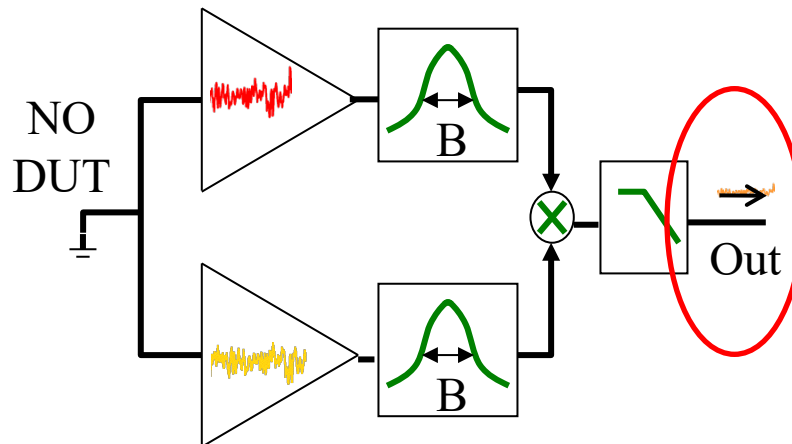
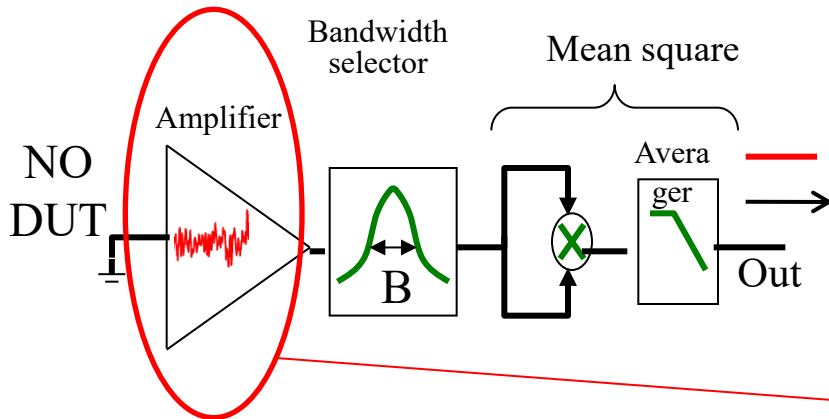
No chance to measure smaller noises ?

A two channels scheme



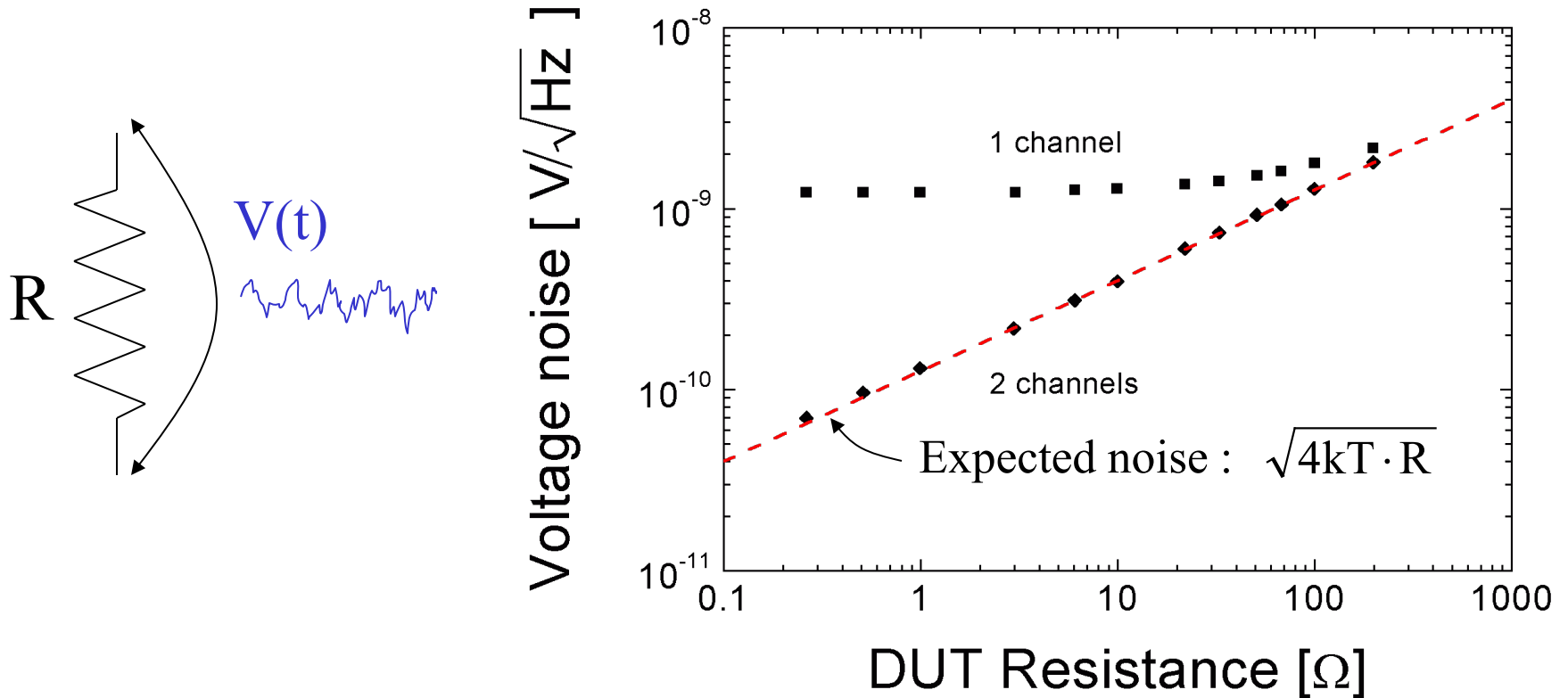
Comparison btwn the two instruments

NO signal applied (NO input DUT)



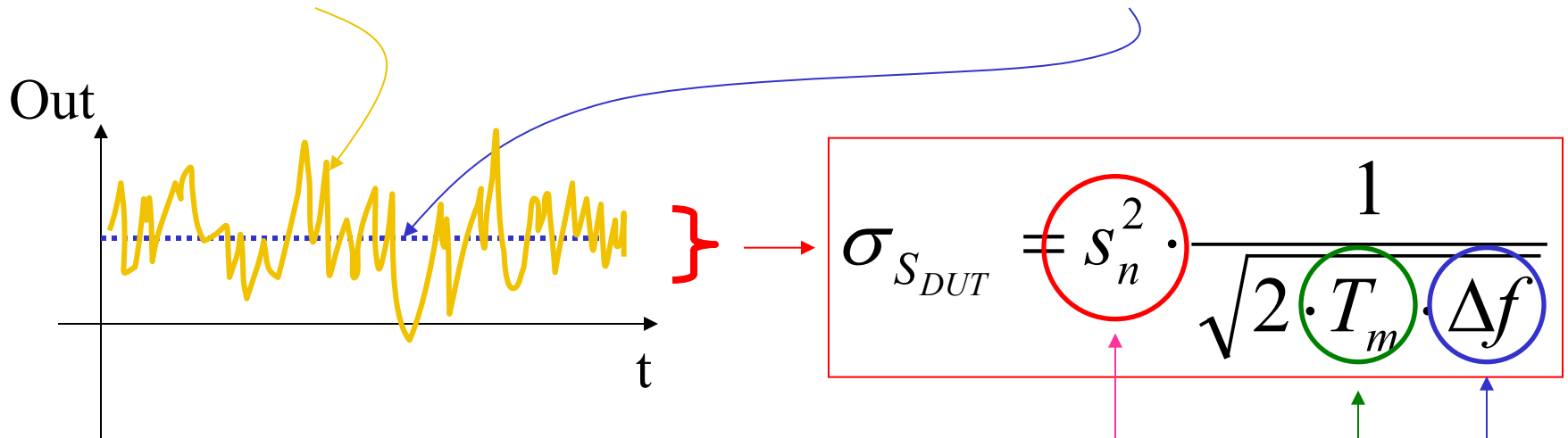
ZERO output (provided sufficiently long averaging)
irrespective to the type of input noise ($1/f$, white, f)

An example



Sensitivity of the instrument

The **level of fluctuation** around the **DUT value** [V^2] is :

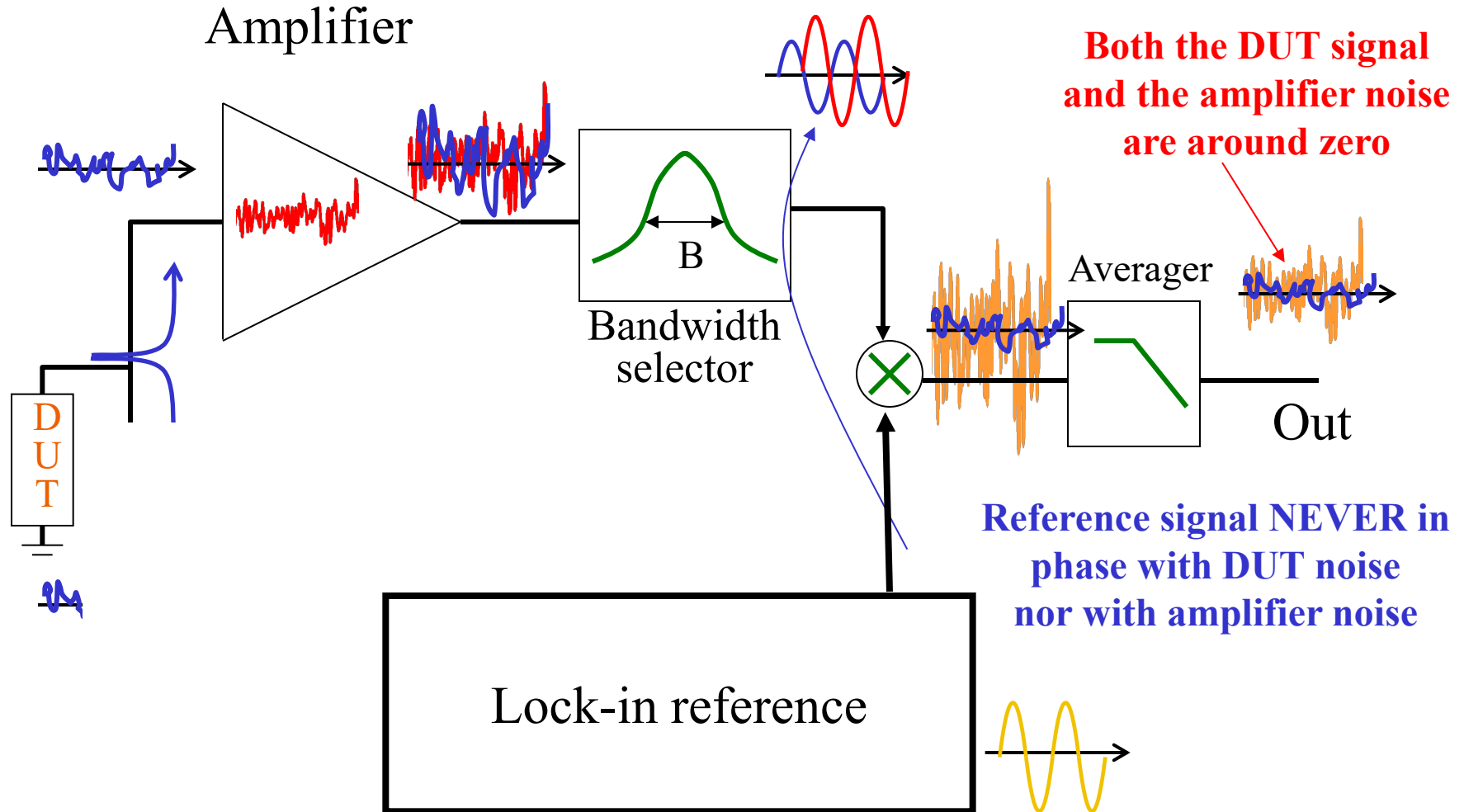


Instrument sensitivity is imposed by:

- the **amplifier input noise**
- the available **time for the measurement**
- the desired **frequency resolution**

Note that σ_{SDUT} [V^2/Hz] or [A^2/Hz] gives the amount of noise power obtained with T_m and/or Δf . The noise level in [$V/\sqrt{\text{Hz}}$] or [$A/\sqrt{\text{Hz}}$] is obtained as $\sqrt{\sigma_{SDUT}}$.

NOT possible with a standard lock-in

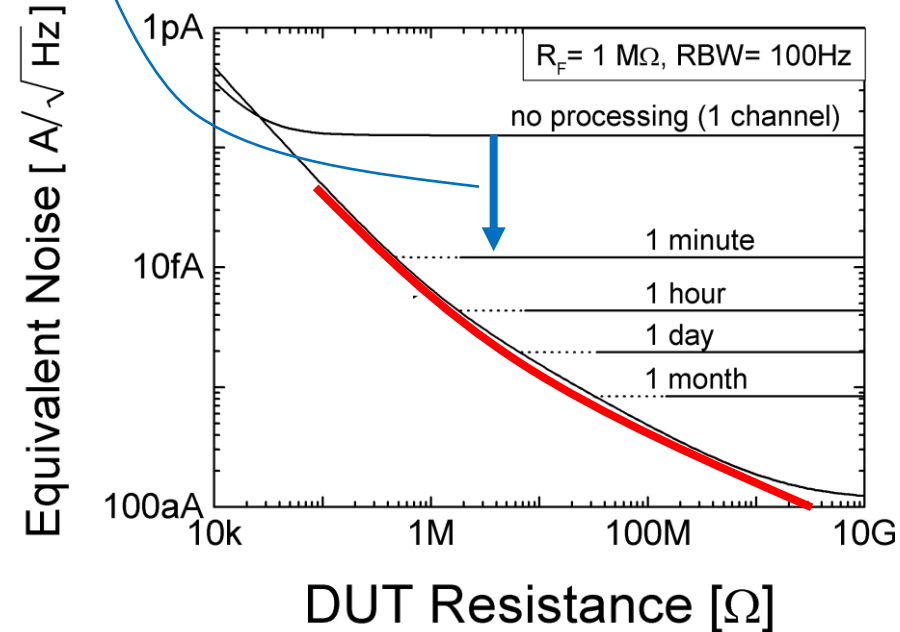
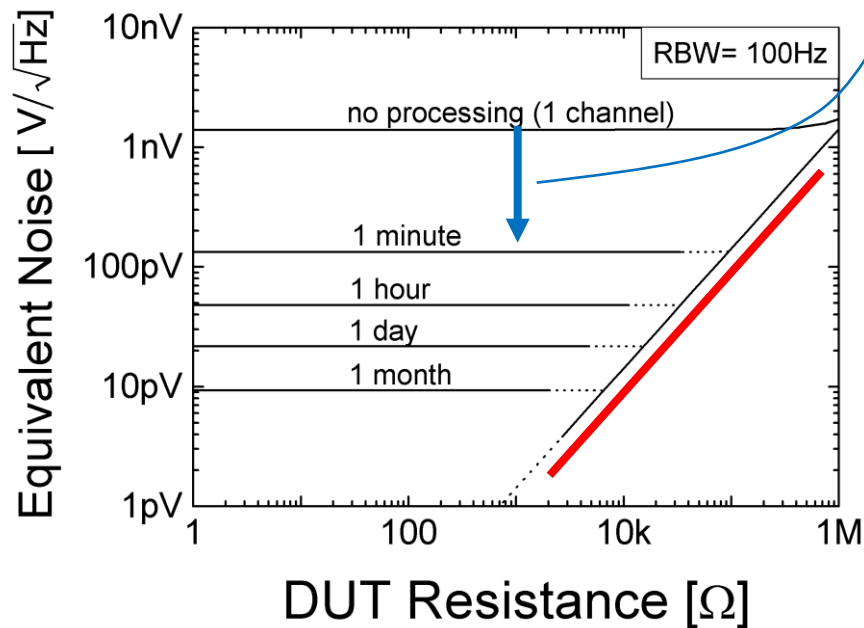


Sensitivity limit vs measuring time

$$\sqrt{S_n^2 \cdot \frac{1}{\sqrt{2 \cdot T_m \cdot \Delta f}}}$$

Voltage noise

Current noise



Limits imposed by residual correlations (see additional material)

Example of current measurement

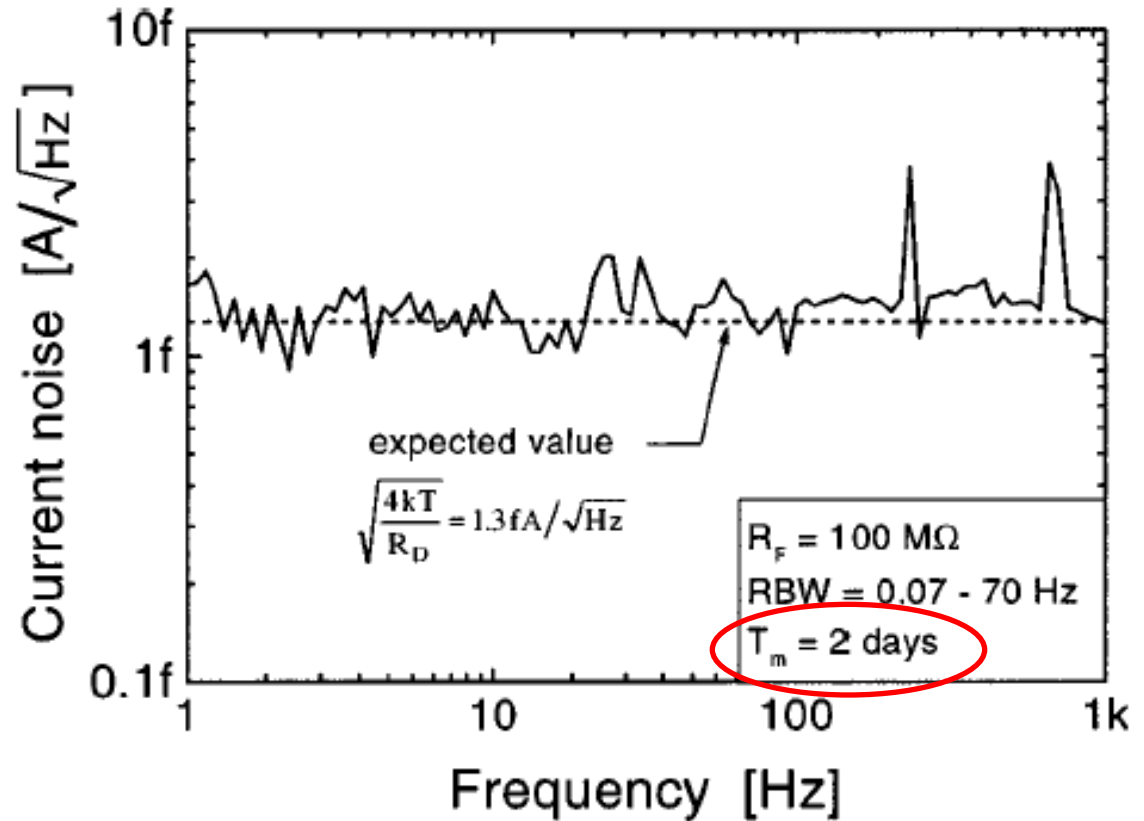
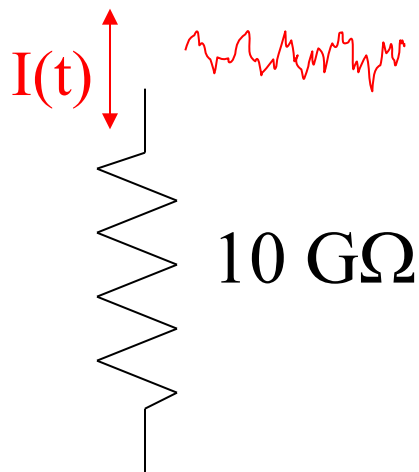
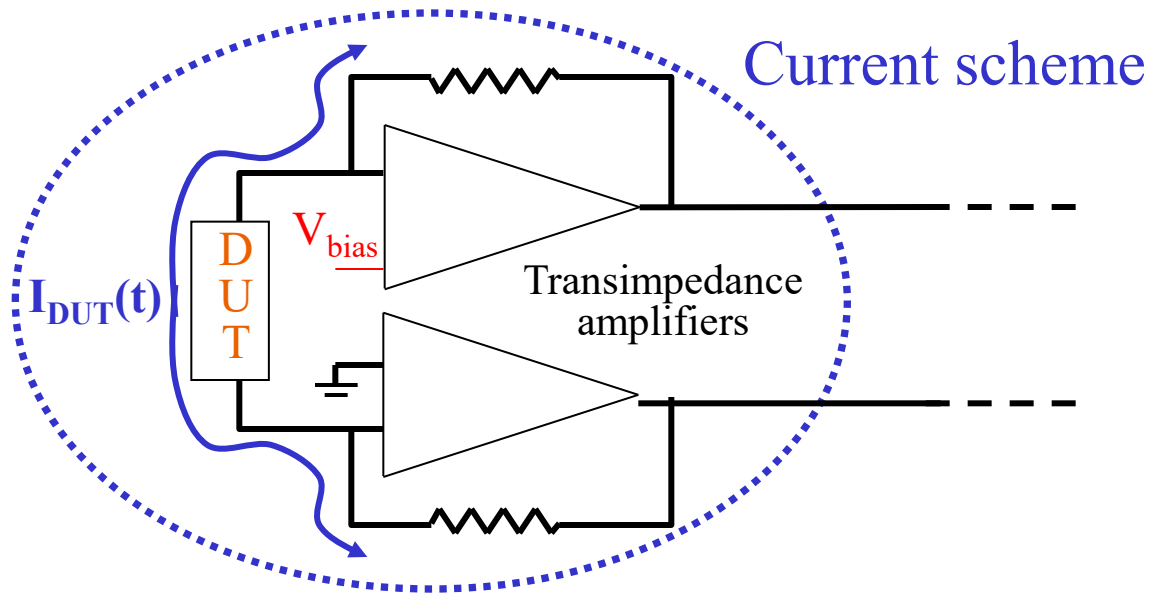
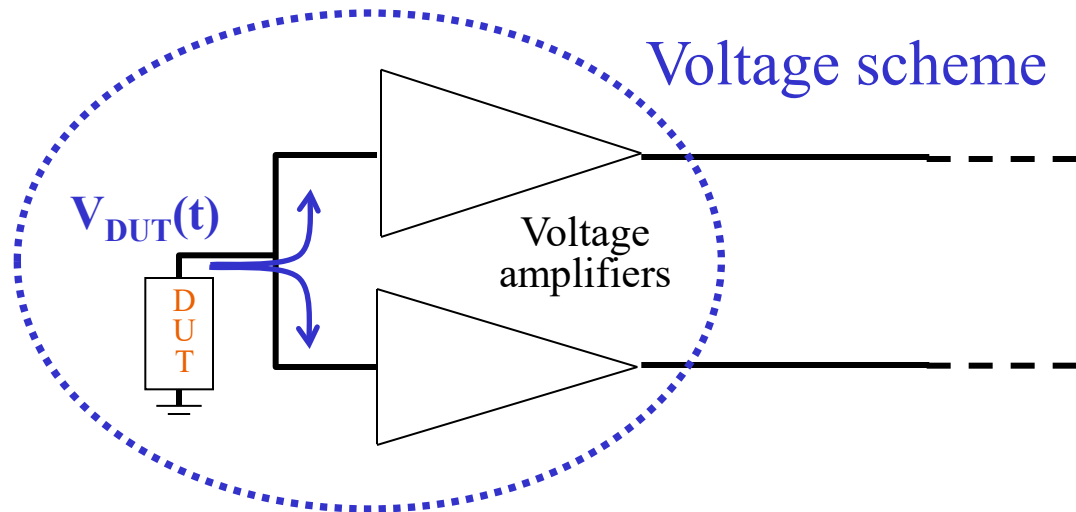


FIG. 7. Frequency spectrum of the current noise produced by a resistor of $10\text{ G}\Omega$. Peaks are probably due to an imperfect shielding from interferences that produce correlated signals.



Current vs Voltage scheme

Choice dictated by DUT impedance to minimize residual correlations:

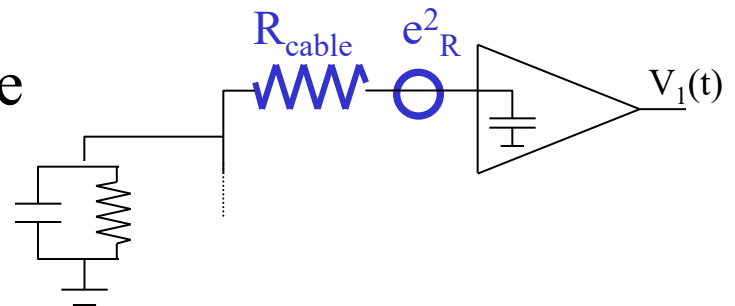
- low impedance DUT ($< 10 \text{ k}\Omega$) \rightarrow voltage mode
- high impedance DUT ($> 10 \text{ k}\Omega$) \rightarrow current mode

Current scheme has practical advantages :

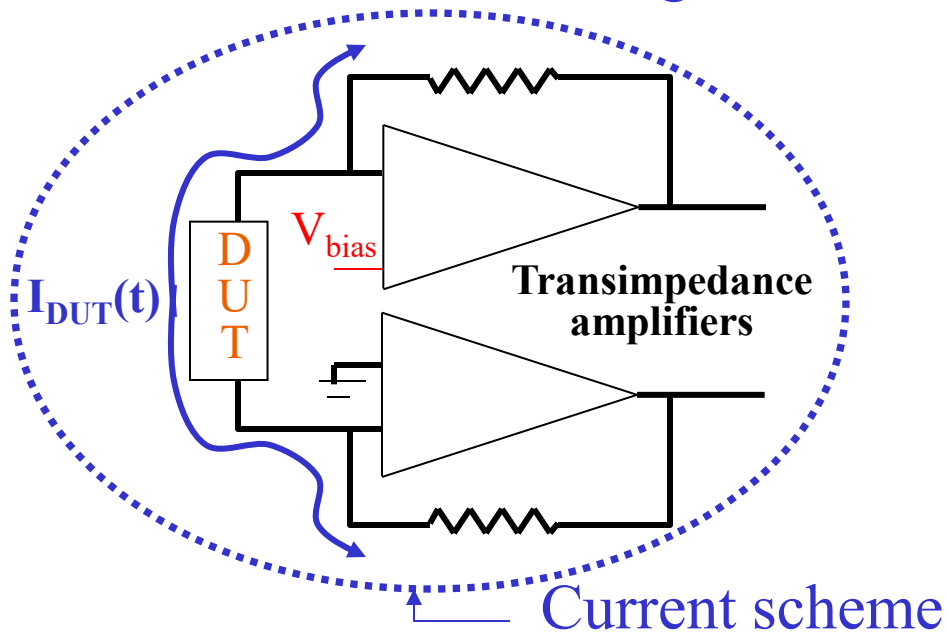
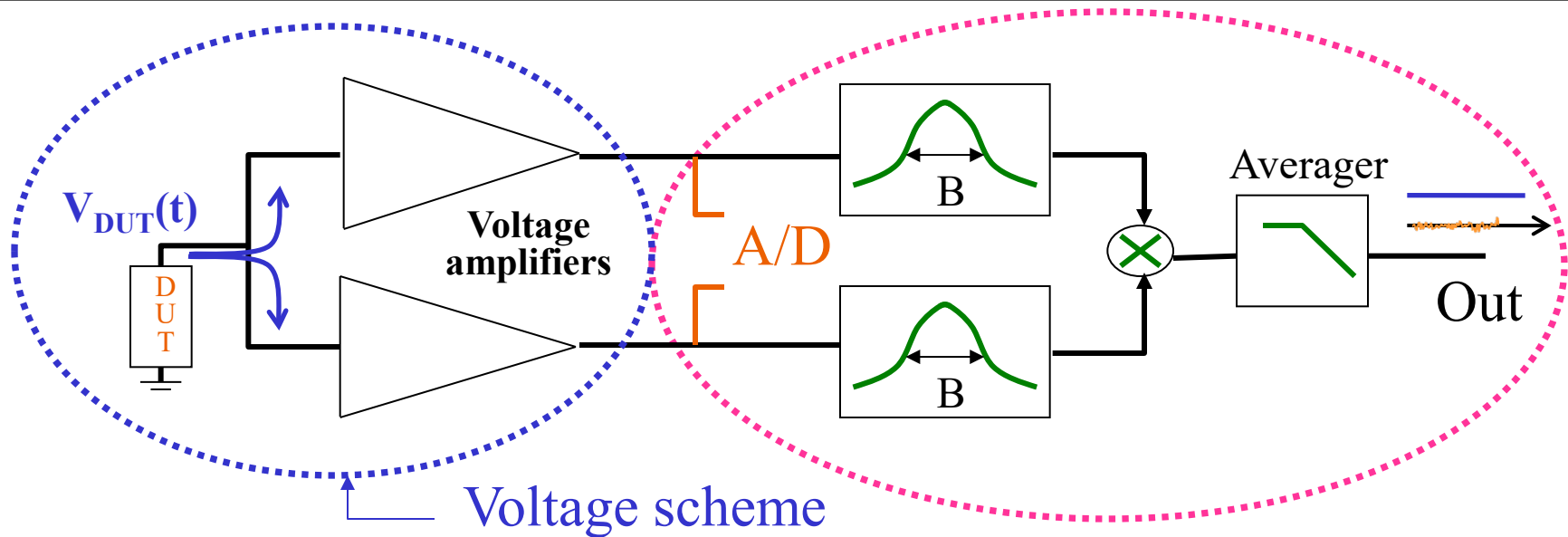
- DUT can be bias directly by the instrument
- DUT biasing network produces less correlated noise

Voltage scheme :

- Noise from cables are less effective



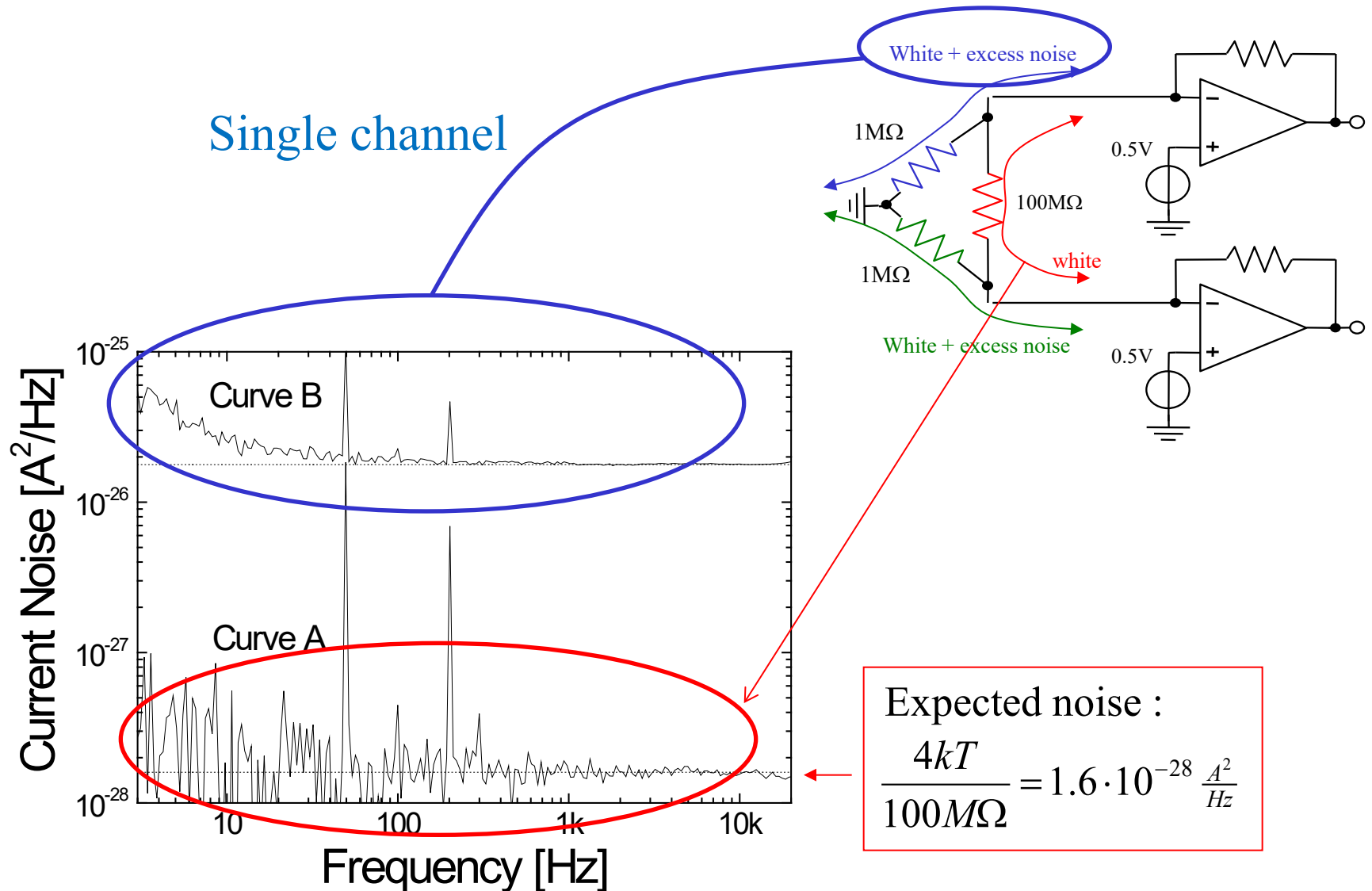
Practical realisation of the instrument



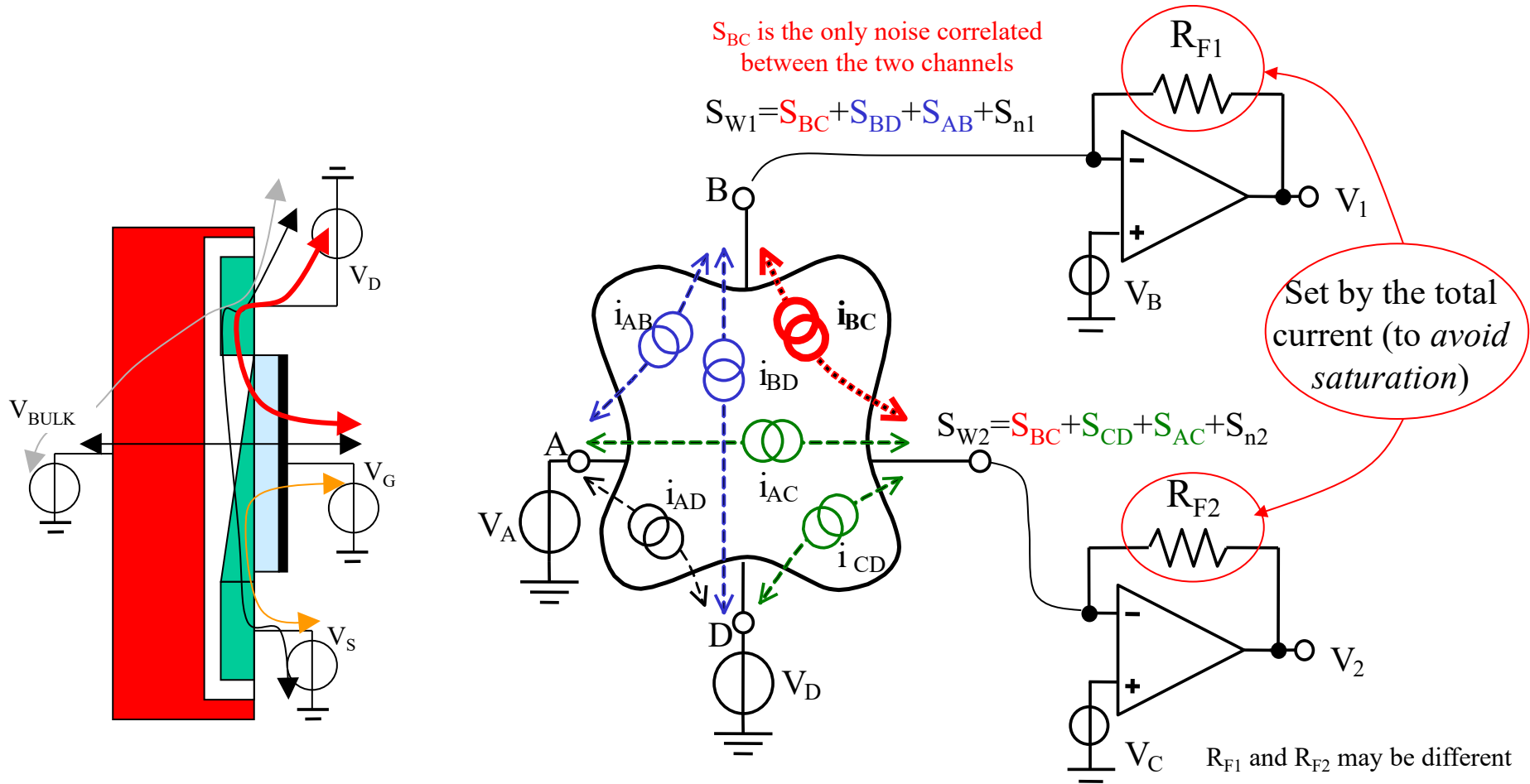
Analog processing :
1 point of the spectrum
acquired at a time
(needs N measurements)

Digital processing :
ALL points of the spectrum
acquired at the same time
(eq. to N filters in parallel)

Extraction of single noise from multipoles



Extraction of single noise from multipoles



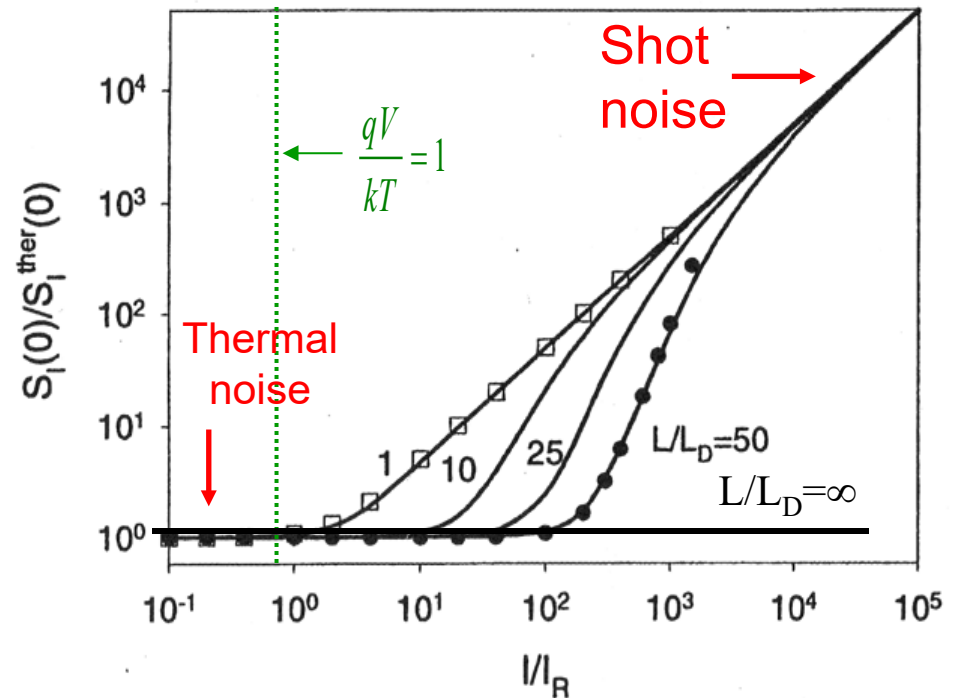
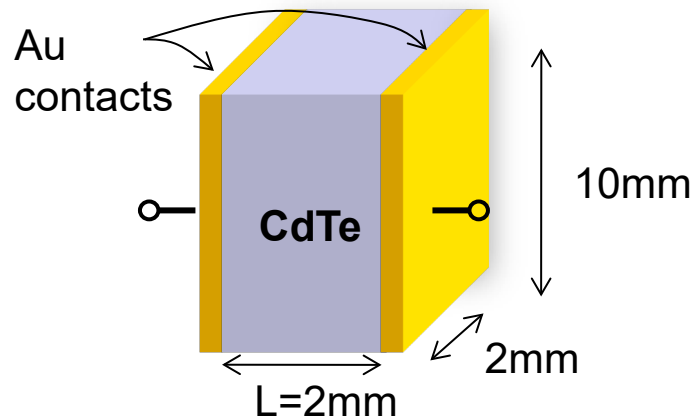
$$T_m \approx \frac{1}{2 \cdot BW} \frac{S_{W1} \cdot S_{W2}}{S_{BC}^2}$$

$$\left. \begin{array}{l} Ex : S_{W1} = S_{W2} = 100 S_{BC} \\ BW = 100 \text{ Hz} \end{array} \right\} T_m \sim 2 \text{ min}$$

In summary ...

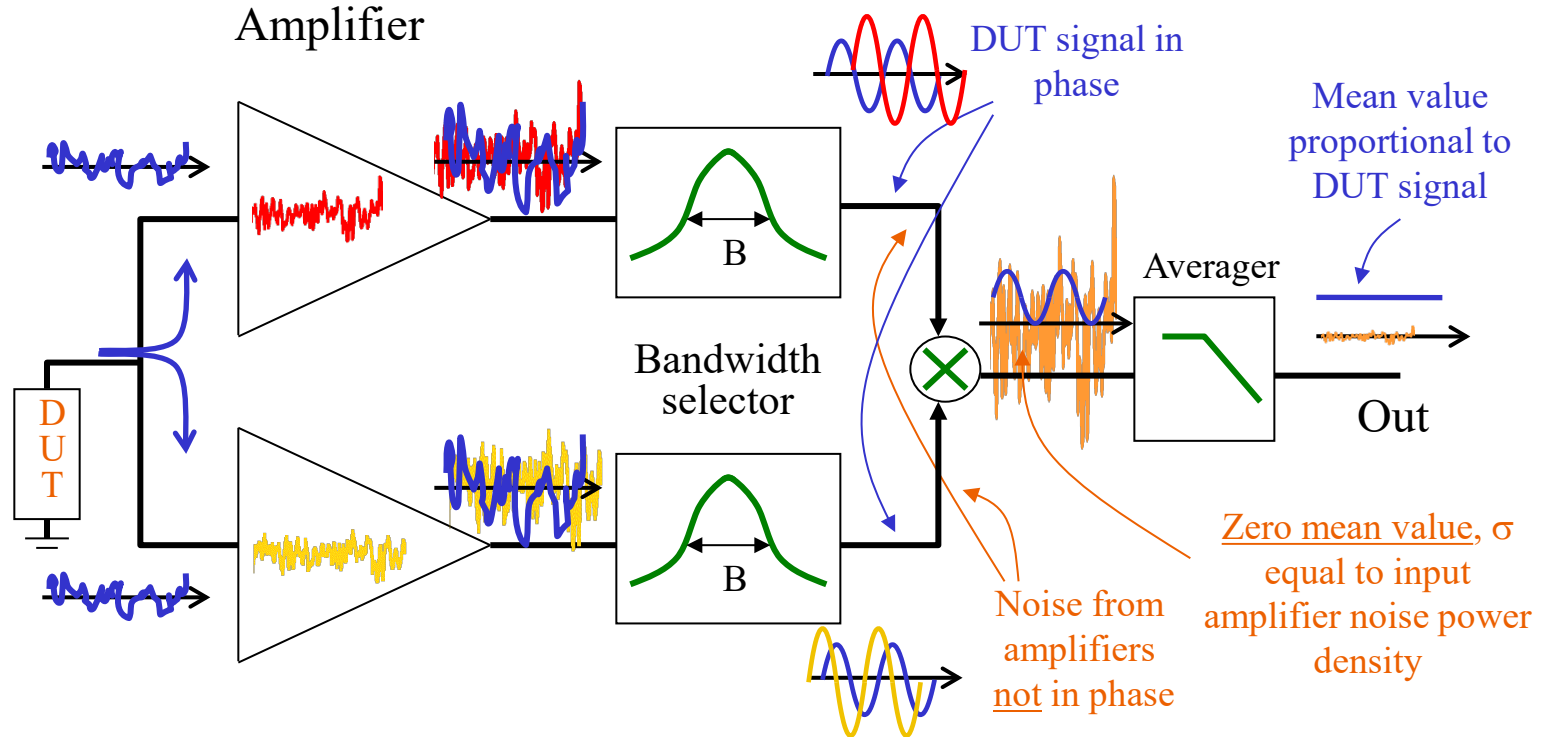
Things to remember (1)

Noise can be an interesting «signal»



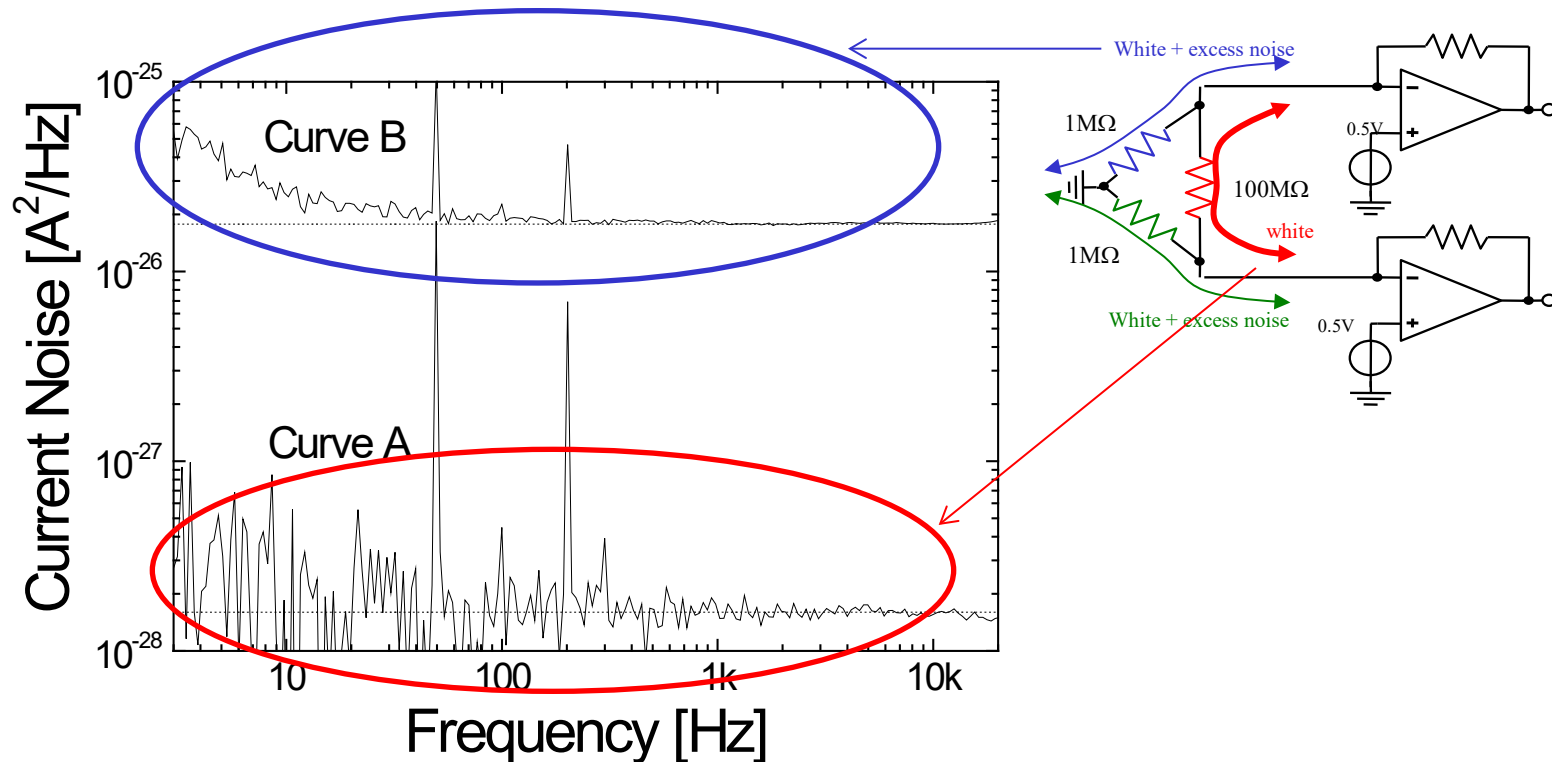
Things to remember (2)

Very small noise can be measured by using a **Correlation Spectrum Technique**



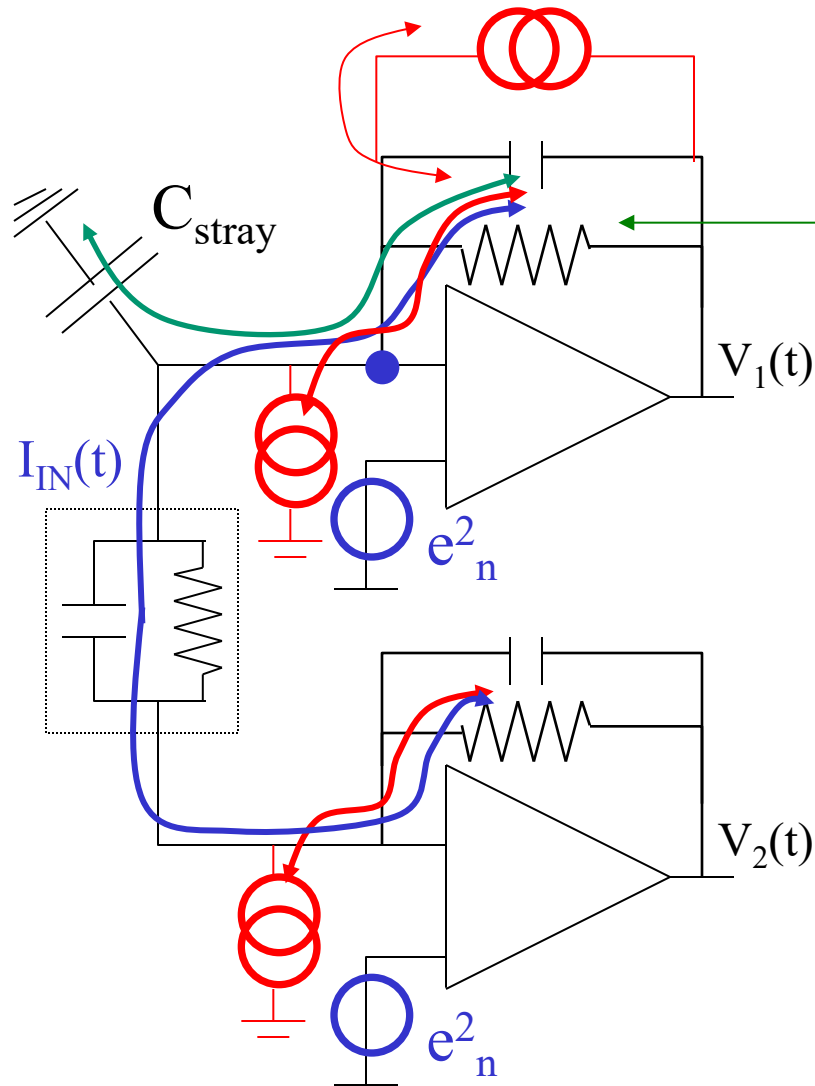
Things to remember (3)

Sorting a noise among many others is possible



ADDITIONAL MATERIAL

Residual correlations : current scheme

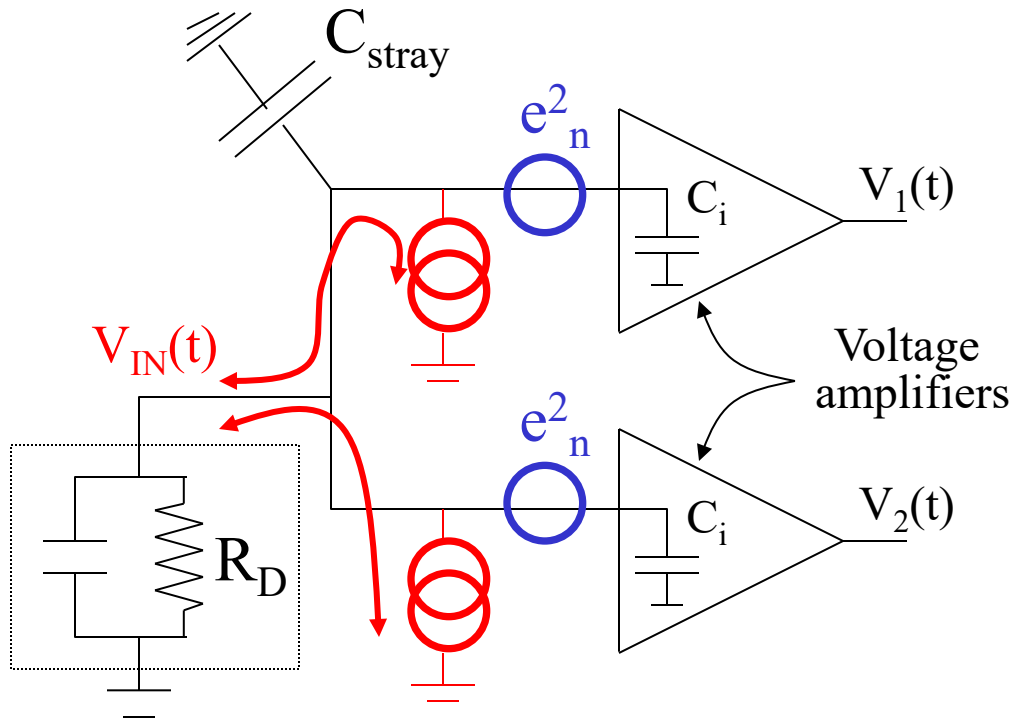


R_F is limited by:

- dynamic range of $I_{DUT|DC}$
- bandwidth

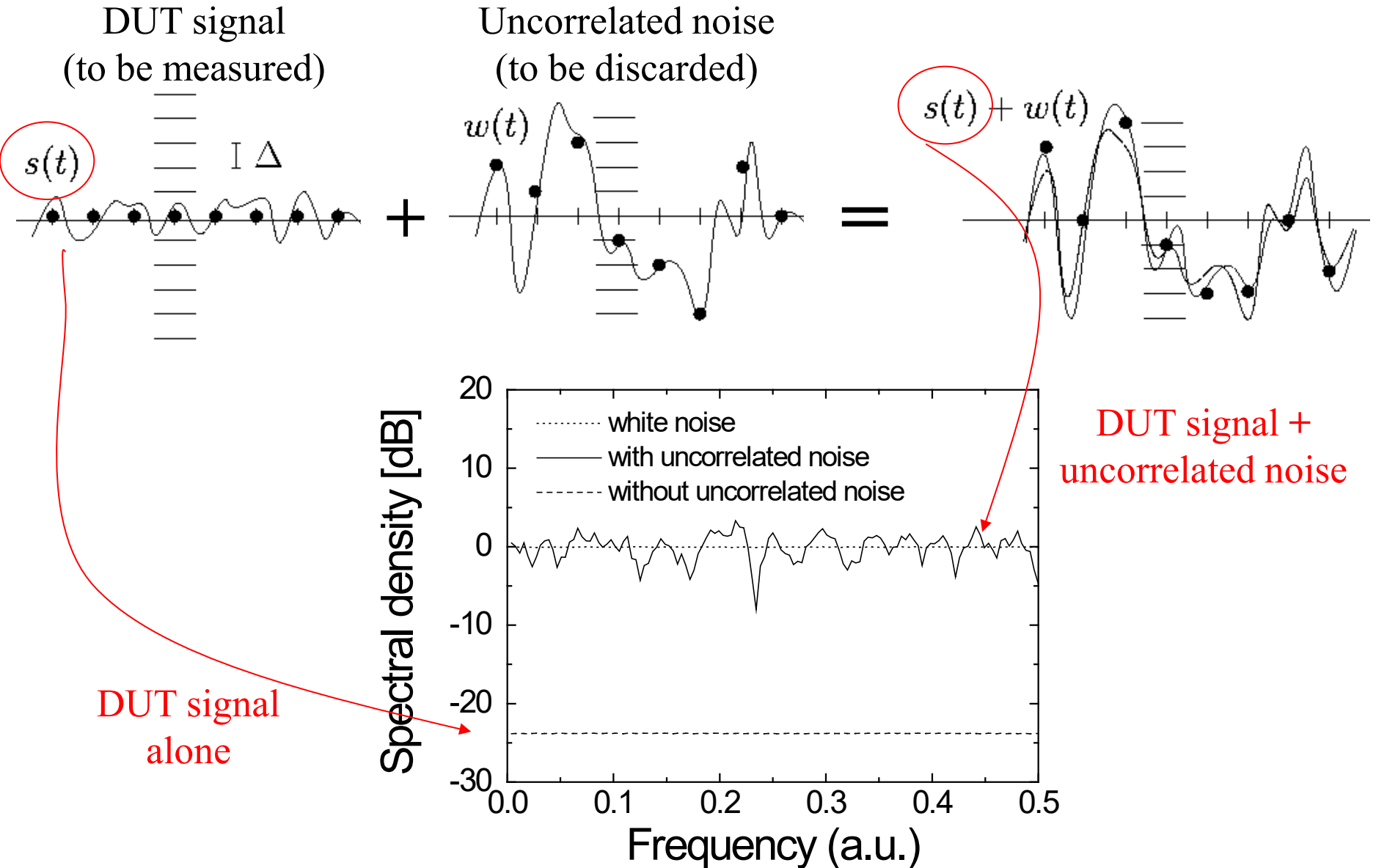
- **Current noises of TIA:**
uncorrelated
(is reduced by averaging)
- **Voltage noise of TIA:**
partly correlated
(is not reduced and sets the sensitivity limit)
partly uncorrelated
(is reduced by averaging)

Residual correlations : voltage scheme



- Current noise : **correlated**
- Voltage noise : **uncorrelated**

Suppression of the quantisation noise



Bandwidth limitations

~ mHz

- Very long measurement time (stationary noise over a long period)
- High stability of the DUT, of the amplifiers and of the connections
- Stability in the set-up temperature

~ MHz

- Set by residual correlation through the DUT ($C_D=10\text{pF}$) and by parasitics ($C_i=10\text{pF}$)
- and by the low noise amplifier bandwidth.
- Needs high sampling frequency in digital mode

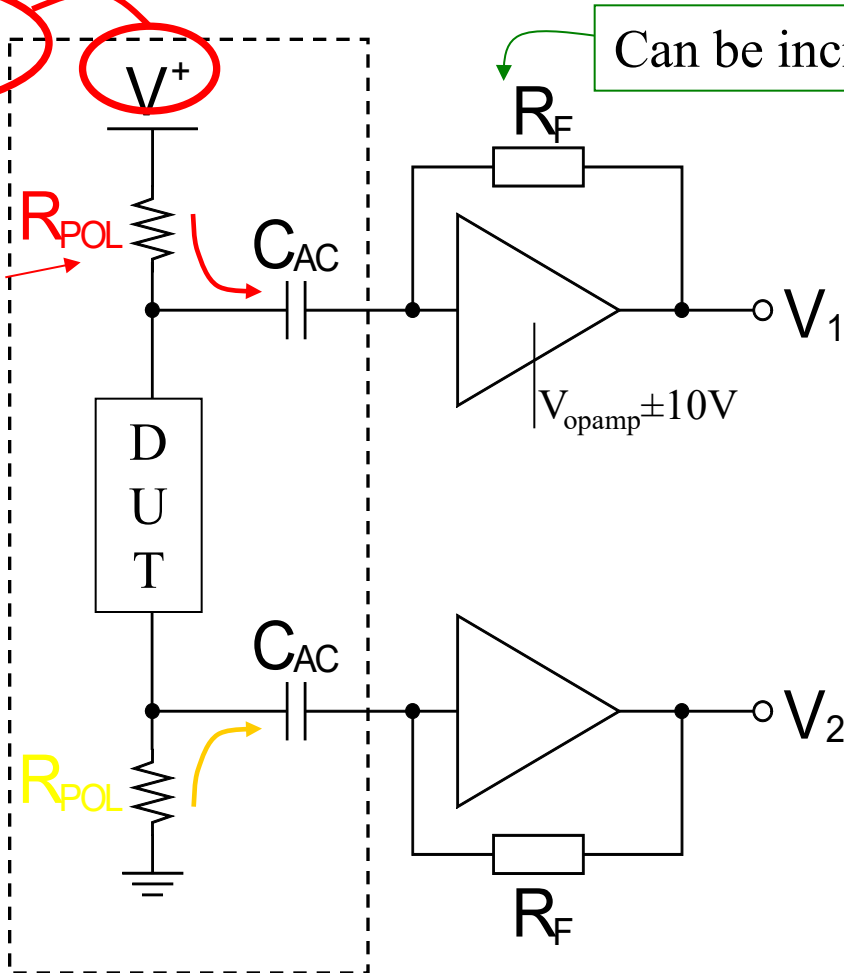
~ GHz

- Reduction of C_D and parasitics (integrated DUT and amplifiers)
- High bandwidth amplifiers (voltage scheme favored).
- Heterodyne system

AC coupling of the DUT – current scheme

Large bias

The largest possible
($V^+ \gg V_{\text{opamp}}$)



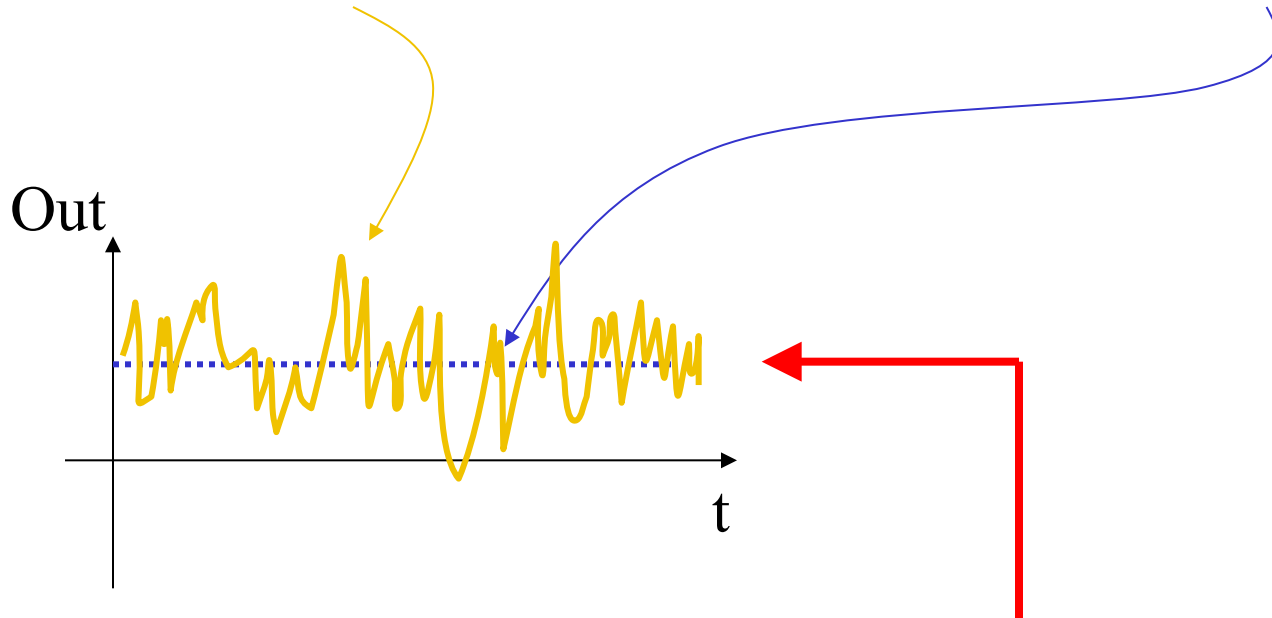
Can be increased (no I_{DC}) \rightarrow reduced T_m

In a *current scheme* :
noise of bias elements
are **uncorrelated**

Example: $V^+ = 100\text{V}$ instead of $V_{\text{opamp}} = 10 \rightarrow S_R(f)/5 \rightarrow T_m/25$!

Accuracy of the instrument

level of fluctuation around the DUT value



Instrument accuracy is limited by the precision in the calculation of :

- the system gain
- the system frequency response